

Hadron Spectroscopy and Dynamics from Light-Front Holography and Conformal Symmetry

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Abstract

To a first semiclassical approximation one can reduce the multi-parton light-front problem in QCD to an effective one-dimensional quantum field theory, which encodes the fundamental conformal symmetry of the classical QCD Lagrangian. This procedure leads to a relativistic light-front wave equation for arbitrary spin which incorporates essential spectroscopic and non-perturbative dynamical features of hadron physics. The mass scale for confinement and higher dimensional holographic mapping to AdS space are also emergent properties of this framework.

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1 Introduction

Recent analytical insights into the nonperturbative nature of the confining interaction in QCD follow from the remarkable holographic correspondence between the equations of motion in AdS space and the light-front (LF) Hamiltonian equations of motion for relativistic light hadron bound-states in physical space-time [1]. In fact, the mapping of the equations of motion [1] and the matching of the electromagnetic [2, 3] and gravitational [4] form factors in AdS space [5, 6] with the corresponding expressions derived from LF quantization in physical space time is the central feature of the LF holographic approach to hadronic physics. This approach allows us to establish a precise relation between wave functions in AdS space and the LF wavefunctions (LFWFs) describing the internal structure of hadrons. However the actual form of the effective potential has remained unknown until very recently and is thus model-dependent.

It was been realized very recently [7] that the form of the effective LF confining potential can be obtained from the framework introduced in a remarkable paper by V. de Alfaro, S. Fubini and G. Furlan (dAFF) [8]. It was shown by dAFF that in the Schrödinger representation, a scale can appear in the Hamiltonian operator while retaining the conformal invariance of the action [8]. This remarkable result is based on the isomorphism of the algebra of the one-dimensional conformal group $\text{Conf}(R^1)$ to the algebra of generators of the group $\text{SO}(2,1)$. One of the generators of this group, the rotation in the 2-dimensional space, is compact and has therefore a discrete spectrum with normalizable eigenfunctions. As a result, the form of the evolution operator is fixed and includes a confining harmonic oscillator potential, and the time variable has a finite range. In fact, it was shown in Ref. [7] that there exists a remarkable holographic connection between the one-dimensional semi-classical approximation to light-front dynamics with gravity in a higher dimensional AdS space, and the constraints imposed by the invariance properties under the full conformal group in one dimension $\text{Conf}(R^1)$. Other approaches to emergent holography are discussed in [9, 10, 11, 12].

2 Light-Front Holography and Conformal Quantum Mechanics

For a hadron with four-momentum P^μ , the generators $P^\mu = (P^-, P^+, \vec{P}_\perp)$, $P^\pm = P^0 \pm P^3$, are constructed canonically from the QCD Lagrangian by quantizing the system on the light-front at fixed LF time x^+ , $x^\pm = x^0 \pm x^3$ [13]. The LF Hamiltonian P^- generates the LF

time evolution $P^-|\phi\rangle = i\frac{\partial}{\partial x^+}|\phi\rangle$, whereas the LF longitudinal P^+ and transverse momentum \vec{P}_\perp are kinematical generators. In the limit of zero quark masses the longitudinal modes decouple from the invariant LF Hamiltonian equation $H_{LF}|\phi\rangle = M^2|\phi\rangle$ with $H_{LF} = P_\mu P^\mu = P^-P^+ - \vec{P}_\perp^2$. We obtain the wave equation [1]

$$\left(-\frac{d^2}{d\zeta^2} - \frac{1-4L^2}{4\zeta^2} + U(\zeta, J)\right)\phi_{n,J,L} = M^2\phi_{n,J,L}, \quad (1)$$

a relativistic single-variable LF Schrödinger equation, where the effective potential U acts on the valence sector of the theory. The effective potential follows from the systematic expression of the higher Fock components as functionals of the lower ones [14]. The variable z of AdS space is identified with the LF boost-invariant transverse-impact variable ζ [2], thus giving the holographic variable a precise definition in LF QCD [1, 2]. For a two-parton bound state $\zeta^2 = x(1-x)b_\perp^2$, where x is the longitudinal momentum fraction and b_\perp is the transverse-impact distance between the quark and antiquark.

Recently we have derived wave equations for hadrons with arbitrary spin starting from an effective action in AdS space [15]. An essential element is the mapping of the higher-dimensional equations to the LF Hamiltonian equation found in Ref. [1]. This procedure allows a clear distinction between the kinematical and dynamical aspects of the LF holographic approach to hadron physics. Accordingly, the non-trivial geometry of pure AdS space encodes the kinematics, and the additional deformations of AdS encode the dynamics, including confinement [15], and determine the form of the LF effective potential from the precise holographic mapping to light-front physics [1, 15]. For $d = 4$ one finds from the dilaton-modified AdS action the LF potential [15, 16]

$$U(\zeta, J) = \frac{1}{2}\varphi''(\zeta) + \frac{1}{4}\varphi'(\zeta)^2 + \frac{2J-3}{2\zeta}\varphi'(\zeta), \quad (2)$$

provided that the product of the AdS mass m and the AdS curvature radius R are related to the total and orbital angular momentum, J and L respectively, according to $(mR)^2 = -(2-J)^2 + L^2$. The critical value $J = L = 0$ corresponds to the lowest possible stable solution, the ground state of the LF Hamiltonian, in agreement with the AdS stability bound $(mR)^2 \geq -4$ [17], where R is the AdS radius.

The classical Lagrangian of QCD is, in the limit of massless quarks, invariant under conformal transformations. Since we are interested in a semiclassical approximation to the nonperturbative domain of QCD, analogous to the quantum mechanical wave equations in atomic physics, it is natural to have a closer look at conformal quantum mechanics, a conformal field theory in one dimension, following dAFF [8]. While leaving the action

invariant, the dAFF procedure leads to a redefinition of the quantum mechanical evolution operator, and consequently to a redefinition of the corresponding ‘time’ evolution parameter τ , the range of which is finite. In the Schrödinger representation

$$i\frac{\partial}{\partial\tau}\psi(x,\tau) = H_\tau\left(x, -i\frac{d}{dx}\right)\psi(x,\tau),$$

the dAFF Hamiltonian H_τ is given by [8, 7]

$$H_\tau = \frac{1}{2}u\left(-\frac{d^2}{dx^2} + \frac{g}{x^2}\right) + \frac{i}{4}v\left(x\frac{d}{dx} + \frac{d}{dx}x\right) + \frac{1}{2}wx^2, \quad (3)$$

which is at $\tau = 0$, the superposition of the ‘free’ Hamiltonian H , the generator of dilatations D and the generator of special conformal transformations K in one dimension, the generators of $\text{Conf}(R^1)$; namely

$$H_\tau = uH + vD + wK.$$

The conformal group $\text{Conf}(R^1)$ is locally isomorphic to $\text{SO}(2,1)$, the Lorentz group in 2+1 dimensions. Since the generators of $\text{Conf}(R^1)$ have different dimensions, their relations with the generators of $\text{SO}(2,1)$ imply a scale, which here plays a fundamental role, as already conjectured in [8].

Comparing the dAFF Hamiltonian (3) with the light-front wave equation (1) and identifying the variable x with the light-front invariant variable ζ , we have to choose $u = 2$, $v = 0$ and relate the dimensionless constant g to the LF orbital angular momentum, $g = L^2 - 1/4$, in order to reproduce the light-front kinematics. Furthermore $w = 2\lambda^2$ fixes the confining light-front potential to a quadratic $\lambda^2\zeta^2$ dependence. The choice of the dilaton profile $\varphi(z) = \lambda z^2$ introduced in [18] thus follows from the requirements of conformal invariance. This specific form for $\varphi(z)$ leads through (2) to the effective LF potential

$$U(\zeta, J) = \lambda^2\zeta^2 + 2\lambda(J - 1),$$

and corresponds to a transverse oscillator in the light-front. The term $\lambda^2\zeta^2$ is determined uniquely by the underlying conformal invariance of classical QCD incorporated in the one-dimensional effective theory, and the constant term $2\lambda(J-1)$ by the embedding space [15, 16].

2.1 A Light-Front Holographic Model for Mesons

From (1) one obtains for the effective LF potential $U(\zeta, J) = \lambda^2\zeta^2 + 2\lambda(J - 1)$ for $\lambda > 0$ a mass spectrum for mesons characterized by the total angular momentum J , the orbital angular momentum L and orbital excitation n given by

$$M_{n,J,L}^2 = 4\lambda\left(n + \frac{J+L}{2}\right),$$

an important result also found in Ref. [19]. This result not only implies linear Regge trajectories, but also a massless pion and the relation between the ρ and a_1 mass usually obtained from the Weinberg sum rules [20]. The model also predicts hadronic LFWFs which underlie form factors [21] and other dynamical observables, as well as vector meson electroproduction [22].

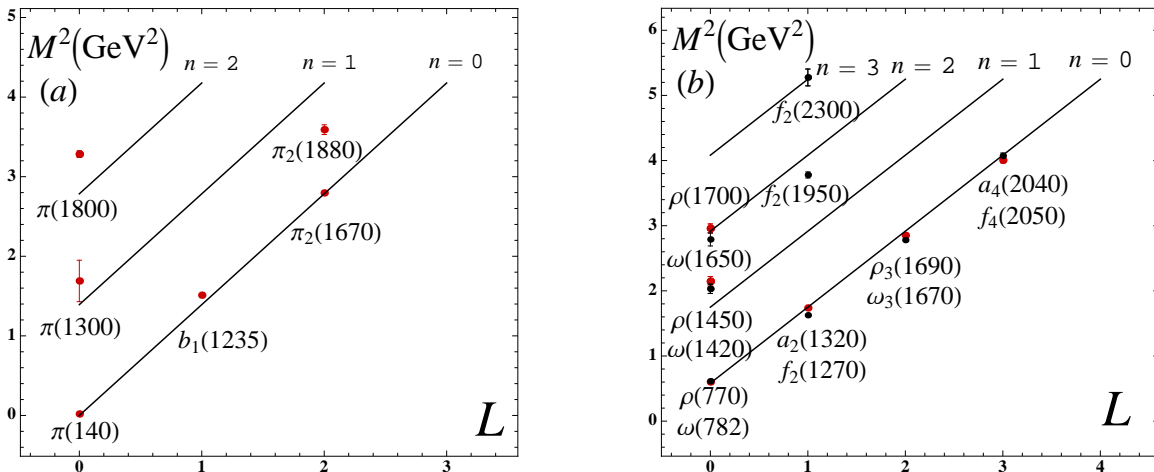


Figure 1: $I=1$ parent and daughter Regge trajectories for the light pseudoscalar mesons (a) with $\sqrt{\lambda} = 0.59$ GeV; and $I=0$ and $I=1$ light vector-mesons (b) with $\sqrt{\lambda} = 0.54$ GeV.

The spectral predictions for the $J = L + S$ light pseudoscalar and vector meson families are compared with experimental data in Fig. 1. The data are from PDG [23]. As we will discuss in the next section, a spin-orbit effect is only predicted for mesons not baryons, as observed in experiment; it thus becomes a crucial test for any model which aims to describe the systematics of the light hadron spectrum. Using the spectral formula for M^2 given above, we find [21]

$$M_{a_2(1320)} > M_{a_1(1260)} > M_{a_0(980)}.$$

The predicted values are 0.76, 1.08 and 1.32 GeV for the masses of the $a_0(980)$, $a_1(1260)$ and $a_2(1320)$, compared with the experimental values 0.98, 1.23 and 1.32 GeV respectively. The prediction for the mass of the $L = 1$, $n = 1$ state $a_0(1450)$ is 1.53 GeV, compared with the observed value 1.47 GeV. The LF holographic model with $\lambda > 0$ accounts for the mass pattern observed in the radial and orbital excitations of the light mesons, as well as for the triplet splitting for the $L = 1$, $J = 0, 1, 2$, vector meson a -states. The slope of the Regge trajectories gives a value $\sqrt{\lambda} \simeq 0.5$ GeV. The value of λ required for describing the pseudoscalar sector is slightly higher than the value of λ extracted from the vector sector. The prediction for the observed spin-orbit splitting for the $L = 1$ a -vector mesons is overestimated

by the model. Note that the solution for $\lambda < 0$ leads to a pion mass heavier than the ρ meson in clear disagreement with observations.

2.2 A Light-Front Holographic Model for Baryons

The analytical exploration of the baryon spectrum using light-front holographic ideas is not as simple or as well understood as the meson case. However, as we shall discuss in this section, even a relatively simple approach provides a framework for a useful description of the baryon spectrum which gives important insights into its systematics. In a chiral spinor component representation, the light-front wave equations for baryons are given by the coupled linear differential equations [15]

$$\begin{aligned} -\frac{d}{d\zeta}\psi_- - \frac{\nu + \frac{1}{2}}{\zeta}\psi_- - V(\zeta)\psi_- &= M\psi_+, \\ \frac{d}{d\zeta}\psi_+ - \frac{\nu + \frac{1}{2}}{\zeta}\psi_+ - V(\zeta)\psi_+ &= M\psi_-, \end{aligned} \quad (4)$$

where ν is given in terms of the mass appearing in the Dirac equation in AdS space: $\nu = |\mu R| - \frac{1}{2}$. One can also show that the effective potential V is J -independent [15, 19]. This is a remarkable result, since it implies that independently of the specific form of the potential, the value of the baryon masses along a given Regge trajectory depends only on the LF orbital angular momentum L , and thus, in contrast with the vector mesons, there is no spin-orbit coupling, in agreement with the observed near-degeneracy in the baryon spectrum.

We choose an effective linear confining potential $V = \lambda \zeta$ which also leads to linear Regge trajectories in both the orbital and radial quantum numbers for baryon excited states. The linear potential also leads to the LF oscillator form $\lambda^2 \zeta^2$ in the second order version of Eqs. 4. For $\lambda > 0$ we find the result

$$M^2 = 4 \lambda (n + \nu + 1).$$

For $\lambda < 0$ no solution is possible [15]. To determine the internal spin, internal orbital angular momentum and radial quantum number assignment of the N and Δ excitation spectrum from the total angular momentum-parity PDG assignment, it is convenient to use the conventional $SU(6) \supset SU(3)_{flavor} \times SU(2)_{spin}$ multiplet structure.

The lowest possible stable state, the proton, corresponds to $n = 0$ and $\nu = 0$. This fixes the scale $\sqrt{\lambda} \simeq 0.5$ GeV. The resulting predictions for the positive-parity spin- $\frac{1}{2}$ nucleons are shown in Fig. 2 (a) for the parent Regge trajectory for $n = 0$ and $\nu = 0, 2, 4, \dots, L$, where L is the relative LF angular momentum between the active quark and the spectator

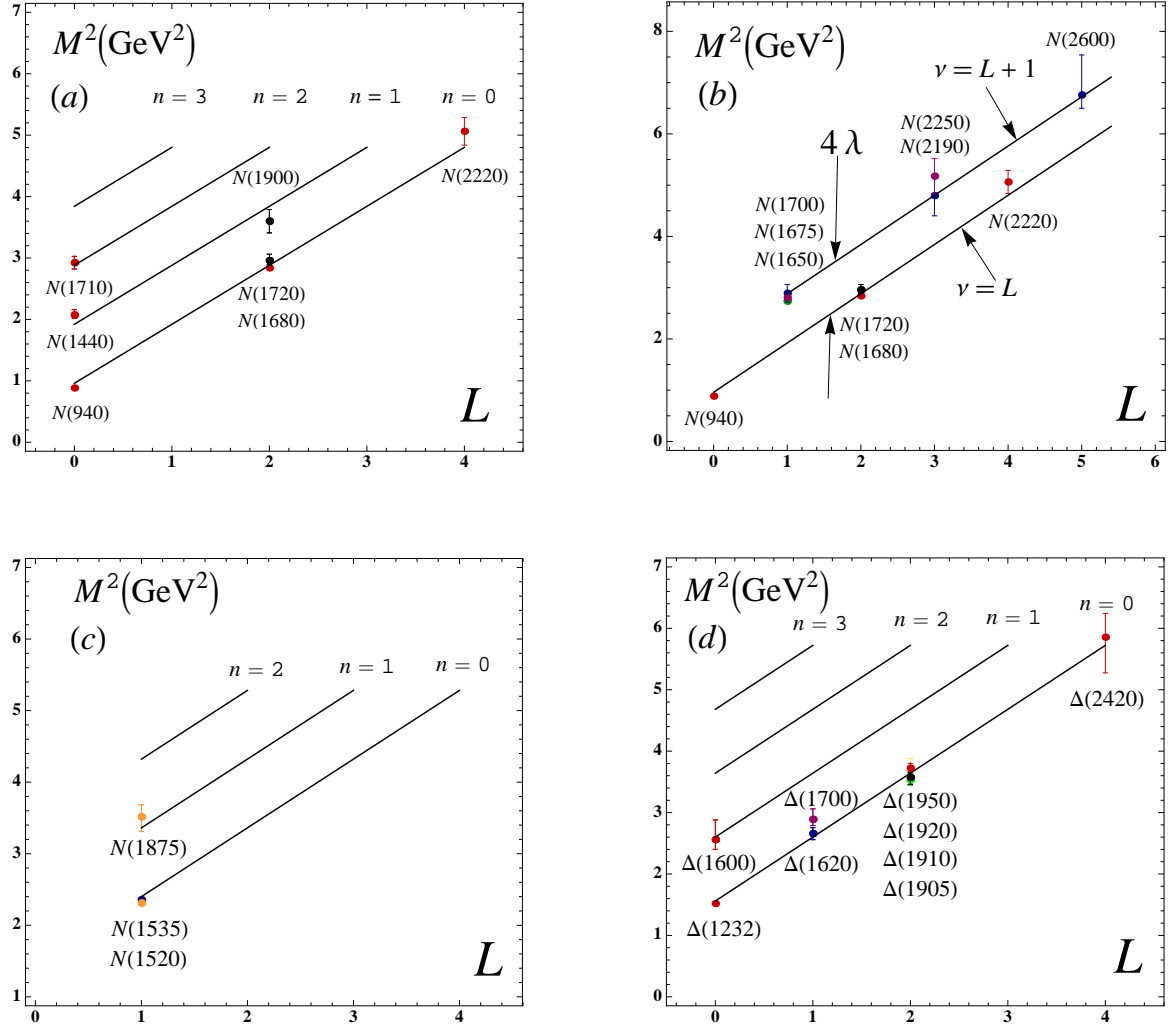


Figure 2: Orbital and radial baryon excitation spectrum. Positive-parity spin- $\frac{1}{2}$ nucleons (a) and spectrum gap between the negative-parity spin- $\frac{3}{2}$ and the positive-parity spin- $\frac{1}{2}$ nucleons families (b). Minus parity N (c) and plus and minus parity Δ families (d), for $\sqrt{\lambda} = 0.49$ GeV (nucleons) and 0.51 GeV (Deltas).

cluster. The predictions for the daughter trajectories for $n = 1, n = 2, \dots$ are also shown in this figure. Only confirmed PDG [23] states are shown. The Roper state $N(1440)$ and the $N(1710)$ are well accounted for as the first and second radial excited states of the proton. The newly identified state, the $N(1900)$ [23] is depicted here as the first radial excitation of the $N(1720)$. The model is successful in explaining the parity degeneracy observed in the light baryon spectrum, such as the $L = 2$, $N(1680) - N(1720)$ pair in Fig. 2 (a). In Fig. 2 (b) we compare the positive parity spin- $\frac{1}{2}$ parent nucleon trajectory with the negative parity

spin- $\frac{3}{2}$ nucleon trajectory. It is remarkable that the gap scale 4λ determines not only the slope of the trajectories, but also the gap in the spectrum between the plus-parity spin- $\frac{1}{2}$ and the minus-parity spin- $\frac{3}{2}$ nucleon families, as indicated by arrows in this figure. This means the respective assignment $\nu = L$ and $\nu = L + 1$ for the lower and upper trajectories in Fig. 2 (b). We also note that the degeneracy of states with the same orbital quantum number L is also well described, as for example the degeneracy of the $L = 1$ states $N(1650)$, $N(1675)$ and $N(1700)$ in Fig. 2 (b).

We have also to take into account baryons with negative parity and internal spin $S = \frac{1}{2}$, as well as baryon states with positive parity and internal spin $S = \frac{3}{2}$ such as the $\Delta(1232)$. Those states are well described by the assignment $\nu = L + \frac{1}{2}$. This means, for example, that $M_{n,L,S=\frac{3}{2}}^{2(+)} = M_{n,L,S=\frac{1}{2}}^{2(-)}$ and consequently the positive and negative-parity Δ states lie in the same trajectory consistent with the experimental results, as depicted in Fig. 2 (d). The newly found state, the $N(1875)$ [23], depicted in Fig. 2 (c) is well described as the first radial excitation of the $N(1520)$, and the near degeneracy of the $N(1520)$ and $N(1535)$ is also well accounted. Likewise, the $\Delta(1600)$ corresponds to the first radial excitation of the $\Delta(1232)$ as shown in Fig. 2 (d). The model explains the important degeneracy of the $L = 2$, $\Delta(1905)$, $\Delta(1910)$, $\Delta(1920)$, $\Delta(1950)$ states which are degenerate within error bars. Our results for the Δ states agree with those of Ref. [24]. “Chiral partners” such as the $N(1535)$ and the $N(940)$ with the same total angular momentum $J = \frac{1}{2}$, but with different orbital angular momentum are non-degenerate from the onset. To recapitulate, the parameter ν has the internal spin S and parity P assignment given in the table below.

Table 1: Orbital assignment for baryon trajectories according to parity and internal spin.

	$S = \frac{1}{2}$	$S = \frac{3}{2}$
P = +	$\nu = L$	$\nu = L + \frac{1}{2}$
P = -	$\nu = L + \frac{1}{2}$	$\nu = L + 1$

The assignment $\nu = L$ for the lowest trajectory, the proton trajectory, is straightforward and follows from the mapping of AdS to light-front physics. The assignment for other spin and parity baryons states given in Table 1 is phenomenological. It is expected that further analysis of the different quark, or quark–diquark, configurations and symmetries of the baryon wave function will indeed explain the actual assignment in Table 1, which successfully describes the full light baryon orbital and radial excitation spectrum, and in particular the gap between trajectories with different parity and internal spin. The holographic variable ζ

has a cluster decomposition and thus labels a system of n -partons as an active quark plus a system of $n - 1$ spectators [2]. From this perspective, a baryon with $n = 3$ looks in light-front holography as a quark–diquark system.

3 Conclusions

We have followed the remarkable results of De Alfaro, Fubini and Furlan [8] which, combined with light-front holographic QCD [1, 7], give important insights into the QCD confining mechanism. It turns out that it is possible to introduce a scale by modifying the variable of dynamical evolution and nonetheless the underlying action stays conformally invariant. This procedure determines uniquely the form of the LF effective potential, corresponding to a quadratic dilaton modification of AdS space to include confinement. We obtain an effective one-dimensional quantum field theory which retains the fundamental conformal symmetry of the classical QCD Lagrangian in the limit of massless quarks. As a result the mass scale for confinement and holographic mapping to AdS space are emergent properties. The group theoretical arguments based on the underlying conformality fix the quadratic form of the effective dilaton profile and thus the corresponding form of the confinement potential of the LF QCD Hamiltonian. The result is a relativistic LF wave equation for bound states which leads to a remarkable description of the spectroscopy and dynamics of the light mesons and baryons [21].

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References

- [1] G. F. de Teramond and S. J. Brodsky, “Light-front holography: a first approximation to QCD,” *Phys. Rev. Lett.* **102**, 081601 (2009) [[arXiv:0809.4899](#) [[hep-ph](#)]].
- [2] S. J. Brodsky and G. F. de Teramond, “Hadronic spectra and light-front wavefunctions in holographic QCD,” *Phys. Rev. Lett.* **96**, 201601 (2006) [[arXiv:hep-ph/0602252](#)].

- [3] S. J. Brodsky and G. F. de Teramond, “Light-front dynamics and AdS/QCD correspondence: the pion form factor in the space- and timelike regions,” *Phys. Rev. D* **77**, 056007 (2008) [[arXiv:0707.3859 \[hep-ph\]](#)].
- [4] S. J. Brodsky and G. F. de Teramond, “Light-front dynamics and AdS/QCD correspondence: gravitational form factors of composite hadrons,” *Phys. Rev. D* **78**, 025032 (2008) [[arXiv:0804.0452 \[hep-ph\]](#)].
- [5] J. Polchinski and M. J. Strassler, “Deep inelastic scattering and gauge/string duality,” *JHEP* **0305**, 012 (2003) [[arXiv:hep-th/0209211](#)].
- [6] Z. Abidin and C. E. Carlson, “Gravitational form factors of vector mesons in an AdS/QCD model,” *Phys. Rev. D* **77**, 095007 (2008) [[arXiv:0801.3839 \[hep-ph\]](#)].
- [7] S. J. Brodsky, G. F. de Teramond and H. G. Dosch, “Threefold complementary approach to holographic QCD,” *Phys. Lett. B* **729**, 3 (2014) [[arXiv:1302.4105 \[hep-th\]](#)].
- [8] V. de Alfaro, S. Fubini and G. Furlan, “Conformal invariance in quantum mechanics,” *Nuovo Cim. A* **34**, 569 (1976).
- [9] R. d. M. Koch, A. Jevicki, K. Jin and J. P. Rodrigues, “AdS₄/CFT₃ Construction from collective fields,” *Phys. Rev. D* **83**, 025006 (2011) [[arXiv:1008.0633 \[hep-th\]](#)].
- [10] S. D. Glazek and A. P. Trawiński, “Model of the AdS/QFT duality,” *Phys. Rev. D* **88**, 105025 (2013) [[arXiv:1307.2059 \[hep-ph\]](#)].
- [11] X. -L. Qi, “Exact holographic mapping and emergent space-time geometry,” [arXiv:1309.6282 \[hep-th\]](#).
- [12] D. D. Dietrich, “Hadrons on the worldline, holography, and Wilson flow,” [arXiv:1312.5718 \[hep-ph\]](#).
- [13] S. J. Brodsky, H. C. Pauli and S. S. Pinsky, “Quantum chromodynamics and other field theories on the light cone,” *Phys. Rept.* **301**, 299 (1998) [[arXiv:hep-ph/9705477](#)].
- [14] H. C. Pauli, “On confinement in a light cone Hamiltonian for QCD,” *Eur. Phys. J. C* **7**, 289 (1999) [[arXiv:hep-th/9809005](#)].
- [15] G. F. de Teramond, H. G. Dosch and S. J. Brodsky, “Kinematical and dynamical aspects of higher-spin bound-state equations in holographic QCD,” *Phys. Rev. D* **87**, 075005 (2013) [[arXiv:1301.1651 \[hep-ph\]](#)].

- [16] G. F. de Teramond and S. J. Brodsky, “Gauge/gravity duality and hadron physics at the light-front,” *AIP Conf. Proc.* **1296**, 128 (2010) [[arXiv:1006.2431](#) [[hep-ph](#)]].
- [17] P. Breitenlohner and D. Z. Freedman, “Stability in gauged extended supergravity,” *Annals Phys.* **144**, 249 (1982).
- [18] A. Karch, E. Katz, D. T. Son and M. A. Stephanov, “Linear confinement and AdS/QCD,” *Phys. Rev. D* **74**, 015005 (2006) [[arXiv:hep-ph/0602229](#)].
- [19] T. Gutsche, V. E. Lyubovitskij, I. Schmidt, A. Vega, “Dilaton in a soft-wall holographic approach to mesons and baryons,” *Phys. Rev. D* **85**, 076003 (2012) [[arXiv:1108.0346](#) [[hep-ph](#)]].
- [20] S. Weinberg, “Precise relations between the spectra of vector and axial vector mesons,” *Phys. Rev. Lett.* **18**, 507 (1967).
- [21] G. F. de Teramond and S. J. Brodsky, “Hadronic form factor models and spectroscopy within the gauge/gravity correspondence,” [arXiv:1203.4025](#) [[hep-ph](#)].
- [22] J. R. Forshaw and R. Sandapen, “An AdS/QCD holographic wavefunction for the ρ meson and diffractive ρ meson electroproduction,” *Phys. Rev. Lett.* **109**, 081601 (2012) [[arXiv:1203.6088](#) [[hep-ph](#)]].
- [23] J. Beringer *et al.* (Particle Data Group), “Review of Particle Physics,” *Phys. Rev. D* **86**, 010001 (2012).
- [24] H. Forkel, M. Beyer, T. Frederico, “Linear square-mass trajectories of radially and orbitally excited hadrons in holographic QCD,” *JHEP* **0707**, 077 (2007) [[arXiv:0705.1857](#) [[hep-ph](#)]].